THE DISCHARGE VARIABILITY OF NEOCORTICAL NEURONS DURING HIGH-CONDUCTANCE STATES

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Abstract—In vivo recordings have shown that the discharge of cortical neurons is often highly variable and can have statistics similar to a Poisson process with a coefficient of variation around unity. To investigate the determinants of this high variability, we analyzed the spontaneous discharge of Hodgkin-Huxley type models of cortical neurons, in which in vivo-like synaptic background activity was modeled by random release events at excitatory and inhibitory synapses. By using compartmental models with active dendrites, or single compartment models with fluctuating conductances and fluctuating currents, we found that a high discharge variability was always paralleled with a high-conductance state, while some active and passive cellular properties had only a minor impact. Furthermore, a balance between excitation and inhibition was not a necessary condition for high discharge variability. We conclude that the fluctuating high-conductance state caused by the ongoing activity in the cortical network in vivo may be viewed as a natural determinant of the highly variable discharge of these neurons. © 2003 IBRO. Published by Elsevier Science Ltd. All rights reserved.

Key words: cerebral cortex, computational models, coefficient of variation, Poisson process, conductance-based models.

Cortical neurons in vivo show a highly irregular discharge activity, both during sensory stimuli (e.g. Dean, 1981; Tolhurst et al., 1983; Softky and Koch, 1993; Holt et al., 1996; Shadlen and Newsome, 1998; Stevens and Zador, 1998; Shinomoto et al., 1999) and during spontaneous activity (e.g. Smith and Smith, 1965; Noda et al., 1970; Burns and Webb, 1976). The coefficient of variation (Cv), the standard measure for the variability of a spike train (see Experimental Procedures), was found to be higher than 0.5 for firing frequencies above 30 Hz in cat and macaque V1 and MT neurons (Softky and Koch, 1993). A Cv of 0.8 was reported as the lower limit under in vivo conditions by investigating the responses of individual MT neurons of alert macaque monkeys driven with constant-motion stimuli (Stevens and Zador, 1998).

Much theoretical work has since been devoted to find neuronal mechanisms responsible for the observed high firing irregularity. However, neither the integration of random excitatory postsynaptic potentials (EPSPs) by a simple leaky integrate-and-fire (IAF) neuron model, nor a more biophysical model of a layer V cell with passive dendrites were able to generate the high Cv observed in vivo (Softky and Koch, 1993). To solve this apparent discrepancy, balanced or “concurrent” inhibition and excitation was proposed as a mechanism producing a discharge activity with Poisson-type variability in IAF models (Shadlen and Newsome, 1994; Usher et al., 1994; Troyer and Miller, 1997; Shadlen and Newsome, 1998; Feng and Brown, 1998, 1999), or in single-compartment Hodgkin-Huxley type models (Bell et al., 1995). Later it was demonstrated that, using a leaky integrator model with partial reset mechanism or physiological gain, Poisson-distributed discharge activity at high frequencies can also be obtained without a fine tuning of inhibitory and excitatory inputs (Troyer and Miller, 1997; Christodoulou and Bugmann, 2000, 2001), indicating a possible role of nonlinear spike-generating dynamics for cortical spike-train statistics (Gutkin and Ermentrout, 1998). Recently, the “noisy” aspect of network dynamics was emphasized as a possible mechanism driving cortical neurons to fire irregularly (Usher et al., 1994; Hansel and Sompolinsky, 1996; Lin et al., 1998; Tiesinga and José, 1999). In this context, it was shown that (temporal) correlation in the inputs can produce a high Cv in the cellular response (Stevens and Zador, 1998; Sakai et al., 1999; Feng and Brown, 2000; Salinas and Sejnowski, 2000; for a review see Salinas and Sejnowski, 2001).

The consensus which emerged from these studies is that neurons operating in an excitable or noise-driven regime are capable of showing highly irregular responses. In this subthreshold regime, the membrane potential is close to spike threshold and action potentials are essentially triggered by fluctuations of the membrane potential. In this framework, the irregularity of the discharge and, thus, the Cv value, can be increased by either bringing the membrane closer to firing threshold (e.g. by balancing the mean of excitatory and inhibitory drive; see e.g. Bell et al., 1995; Shadlen and Newsome, 1998; Feng and Brown, 1998, 1999), or by increasing the noise amplitude (e.g. by correlating noisy synaptic inputs; see e.g. Feng and Brown, 2000; Salinas and Sejnowski, 2001). However, the conditions for the appearance of this subthreshold regime, as well as its dependence on various electrophysiological parameters or the characteristics of the driving inputs, are presently unclear.
The subthreshold regime of cortical neurons in vivo consists in large-amplitude voltage fluctuations (Matsumura et al., 1988; Contreras et al., 1996; Nowak et al., 1997; Paré et al., 1998; Azouz and Gray, 1999; Lampi et al., 1999). These fluctuations are caused by the high levels of spontaneous firing (from 5 to 20 Hz in awake animals; see Hubel, 1959; Evarts, 1964; Steriade, 1978; Matsumura et al., 1988; Holmes and Woody, 1989; Steriade et al., 2001) together with the dense interconnectivity of cortical neurons (several thousands of synaptic inputs per neuron; see Szentagothai, 1965; Cragg, 1967; Gruner et al., 1974; DeFelipe and Fariñas, 1992). In addition, this sustained synaptic activity is also responsible for setting cortical neurons into a high-conductance state (Borg-Graham et al., 1996). Synaptic activity is also responsible for setting cortical neurons to operate in a high-conductance state. To that end, we investigated the spike train statistics of spontaneously discharging non-bursting cortical neurons (from 5 to 20 Hz in awake animals; see Hubel, 1959; Evarts, 1964; Steriade, 1978; Matsumura et al., 1988; Contreras et al., 1996; Nowak et al., 1997; Pare´ et al., 1998; Azouz and Gray, 1999; Lampl et al., 2002).

In this paper we investigated the spike train statistics of spontaneously discharging non-bursting cortical neurons operating in a high-conductance state. To that end, we used both detailed multi-compartment models of morphologically reconstructed neocortical pyramidal neurons with active dendrites and a realistic representation of synaptic background activity, as well as simplified single-compartment models which include a two-conductance representation of background activity. Preliminary results have appeared in a conference paper (Rudolph and Destexhe, 2002).

**EXPERIMENTAL PROCEDURES**

Several types of conductance-based models of cortical neurons were considered.

**Detailed biophysical model**

Simulations were performed using a morphologically reconstructed neocortical pyramidal layer VI neuron of cat parietal cortex (Fig. 1A, Contreras et al., 1997). Passive model parameters were adjusted to intracellular recordings obtained after application of TTX and synaptic blockers (Destexhe and Paré, 1999) and kept constant over all simulations. An intracellular resistivity of \( R_m = 250 \, \Omega \, \text{cm}^2 \), membrane resistivity of \( R_m = 22k \, \Omega \, \text{cm}^2 \) (\( R_m = 50 \, \Omega \, \text{cm}^2 \) in axon), and capacitance of \( C_m = 1 \, \mu F \, \text{cm}^{-2} \) (\( C_m = 0.04 \, \mu F \, \text{cm}^{-2} \) in axon) were used, where \( C_m \) was increased and \( R_m \) decreased by a factor of 1.45 to account for the surface correction due to dendritic spines, assuming that about 45% of the dendritic membrane area is represented by spines (DeFelipe and Fariñas, 1992).

Voltage-dependent conductances were inserted in the soma, dendrites and the axon (sodium current \( I_{Na} \), delayed-rectifier potassium current \( I_{K} \), and voltage-dependent potassium current \( I_{K} \)). All currents were described by Hodgkin-Huxley type models (Hodgkin and Huxley, 1952) with kinetics taken from a model of hippocampal pyramidal cells (Traub and Miles, 1991), adjusted to match voltage-clamp data of cortical pyramidal cells (Huguenard et al., 1988). In the standard model, constant peak conductance densities of 52.3 mS cm\(^{-2}\) (36.1 mS cm\(^{-2}\) in soma, 361 mS cm\(^{-2}\) in axon) for \( I_{Na} \), 10.1 mS cm\(^{-2}\) (7 mS cm\(^{-2}\) in soma, 70 mS cm\(^{-2}\) in axon) for \( I_{K} \), and 0.51 mS cm\(^{-2}\) (0.35 mS cm\(^{-2}\) in soma) for \( I_{K} \) were used. These densities correspond to the values found experimentally in adult hippocampal pyramidal neurons (e.g. Magee and Johnston, 1995). To discuss the impact of membrane excitability, in some simulations the peak conductances of all ion channels were scaled by a common multiplicative factor. A "low" excitability (scale factor of 0.6, corresponding to the lower border of the experimental observed range) and a "high" excitability (scale factor of 1.4) deserved special attention.

To test the effects of active membrane properties, in some cases different sets of voltage-dependent conductances were used. This included Ca\(^2+\)-dependent potassium current (C-current) \( I_{Ca} \) (kinetics in Yamada et al., 1989; constant dendritic conductance density of 1 mS cm\(^{-2}\)), a high-threshold Ca\(^2+\) -current (L-current) \( I_{CaL} \) (kinetics in McCormick and Huguenard, 1992; conductance density of 3 mS cm\(^{-2}\) and 1.5 mS cm\(^{-2}\) for proximal and distal dendrites, respectively) and a persistent sodium current \( I_{NaP} \) (kinetics in French et al., 1990; Huguenard and McCormick, 1992; McCormick and Huguenard, 1992; constant dendritic conductance density of 0.1 mS cm\(^{-2}\)) as well as a model with different kinetics for \( I_{Na} \), \( I_{K} \), and an A-type potassium current \( I_{KA} \) (Migliore et al., 1999).

Synaptic currents were incorporated using two-state kinetic models of glutamate \( \alpha \)-amino-3-hydroxy-5-methyl-4-oxazolopropionic acid (AMPA) and GABA\(_A\) receptor types (Destexhe et al., 1998) with quantal conductances of 869.4 pS (distal region, 600 pS for proximal region, see Fig. 1A) for AMPA and 1200 pS for GABA. These quantal conductances were estimated from miniature synaptic events (Destexhe and Paré, 1999). No metabotropic receptors were included. In some simulations, the quantal conductances were changed by a common factor to investigate the impact of the synaptic background strength. The densities of synapses in different regions of the layer VI cell under consideration were estimated from morphological studies in neocortical pyramidal cells (White, 1989; Larkman, 1991; DeFelipe and Fariñas, 1992), leading to a total of 16,563 glutamatergic and 3376 GABAergic synapses. To perform the simulations in a time-efficient manner, an accelerating algorithm (Lyttton, 1996) was used.

Synaptic background activity was simulated by the random activity of inhibitory and excitatory synapses according to a Poisson process with average rates of \( v_{inh} = 5.5 \, \text{Hz} \) for GABA\(_A\) synapses, and \( v_{exc} = 1.0 \, \text{Hz} \) for AMPA synapses, as estimated from intracellular recordings of pyramidal neurons before and after application of TTX (Paré et al., 1998; Destexhe and Paré, 1999). In some simulations, the firing rates of inhibitory and excitatory synapses were changed. Such changes led to changes in the mean conductances underlying excitation and inhibition and, thus, allowed to shift the average membrane potential relative to firing threshold. Balanced or "concurrent" conditions refer to the situation where the mean inhibitory and excitatory conductances due to synaptic activity are such that the membrane operates close to firing threshold. Independent changes in the firing rate were used to investigate the cellular response in unbalanced conditions, i.e. in situations where the membrane is driven by either pronounced excitation or inhibition.

In addition to the firing rate, the statistics of the synaptic background activity was modified by introducing a correlation in the random background activity. To this end, \( N_C \) independent random release events were redistributed among all \( N \) synapses, which for \( N_C < N \) leads to a co-release of several synapses, whereas every terminal still released according to a Poisson process (see details in Destexhe and Paré, 1999; Rudolph and Destexhe, 2001).
The impact of the cell morphology was addressed by using a simplified three-compartment model which had the same membrane area, input resistance and voltage attenuation as the layer VI cell (Destexhe, 2001, see inset in Fig. 8A). The voltage-dependent currents $I_{Na}$, $I_{Kd}$ and $I_M$ used for the model had densities leading to an average membrane voltage $V_m$ as well as $V_m$ fluc-

Fig. 1. Discharge variability in models of morphologically reconstructed cortical neurons. (A) Morphologically reconstructed neocortical pyramidal layer VI neuron of a cat incorporated in the models. The shaded area indicates the proximal region including all dendritic branches within a radius of 40 $\mu$m from the soma. Inside that region there were no excitatory synapses, whereas inhibitory synapses were spread over the whole dendritic tree. (B) Statistical validity of obtained results. The plots show the $C_v$ as a function of the simulation length for representative examples of the simplified single-compartment model (top) and biophysical model (bottom) at different firing rates. Even for low firing rates (approximately 2 Hz), the $C_v$ values deduced from the spontaneous discharge activity converged already after about 50 s to values obtained from marked extended simulation periods. The grey bars indicate the duration of the simulations used in the present study (see Experimental Procedures). (C) The $C_v$ for spontaneous firing activity as a function of the mean ISI for different levels of background activity. The latter was altered either by a change in the quantal synaptic conductances, or the release frequency at excitatory and inhibitory synaptic terminals (ratio was kept constant) around the standard parameter setup (arrow, see Experimental Procedures). In both cases, correlated background activity with Pearson correlation coefficient of approximately 0.1 was used. The star refers to a spontaneous firing rate of about 15 Hz. Fits: Eq. 5 with $T_R = 22$ mS for changes in frequency and $T_R = 17$ ms for changes in quantal synaptic conductances. (D) Typical ISIH and autocorrelogram (inset) for the standard parameter setup with correlated background activity (Pearson correlation coefficient approximately 0.1) for a spontaneous firing rate of about 15 Hz (indicated by the star in B). Fit: Eq. 6 with $q = 1$, $r = 0.022$ ms$^{-1}$, $\alpha = 4.594$. 

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tions corresponding to the low excitable model described above. The synaptic background activity (number of synapses, quantal conductances, release rates) was the same as in the detailed biophysical model.

**Single-compartment models**

Simplified models of synaptic background activity were constructed by incorporating Ornstein-Uhlenbeck processes (Uhlenbeck and Ornstein, 1930) into a single-compartment model with voltage-dependent currents $I_{Na}$, $I_{K}$, and $I_{m}$ corresponding to the detailed biophysical model. The time-dependent membrane current due to synaptic background activity $I_{syn}(t)$ was described by two different models. First, a fluctuating conductance model (see details in Destexhe et al., 2001), in which $I_{syn}(t)$ was decomposed into two time-dependent conductances (excitatory $g_{e}(t)$ and inhibitory $g_{i}(t)$, each described by one-variable stochastic process similar to the Ornstein-Uhlenbeck process:

$$I_{syn}(t) = g_{e}(t)(V(t) - E_{e}) + g_{i}(t)(V(t) - E_{i})$$ (1)

$$\frac{dg_{e}(t)}{dt} = -\frac{1}{\tau_{e}}(g_{e}(t) - g_{0e}) + \sqrt{D_{e}}W_{e}$$ (2)

where $\tau_{e}$ are average conductances, $\tau_{e}$ are time constants, $D_{e}$ are noise diffusion coefficients, and $W_{e}$ denotes white noise of unit S.D. and zero mean for inhibition and excitation, respectively. White noise is obtained for $\tau_{e} \rightarrow 0$, and $\tau_{e} > 0$ yields "colored" noise. The relation between $D_{e}$ and the S.D. $\sigma_{e}$ of the variables $g_{e}(t)$ is given by $\sigma_{e} = \frac{1}{\sqrt{2\pi}}D_{e}$ (see Gillespie, 1996). Typical parameter values are $E_{e} = 0$ mV, $E_{i} = -75$ mV, $g_{0e} = 0.0121$ $\mu$S, $g_{0i} = 0.0573$ $\mu$S, $\sigma_{e} = 0.025$ $\mu$S, $\sigma_{i} = 2.5\sigma_{e}$, $\tau_{e} = 10.49$ ms, $\tau_{i} = 2.728$ ms.

The second model used was a fluctuating current model, where $I_{syn}(t)$ was directly described by an one-variable Ornstein-Uhlenbeck process without the involvement of conductances:

$$\frac{dI_{syn}(t)}{dt} = -\frac{1}{\tau_{syn}}(I_{syn}(t) - I_{syn0}) + \sqrt{D_{syn}}W_{syn}$$ (3)

where $\sigma_{v} = \frac{1}{\sqrt{2\pi}}D_{syn}$ (see e.g. Tiesinga et al., 2000; Brunel et al., 2001). Typical parameter values are $I_{syn0} = 0$ nA, $\sigma_{v} = 0.4$ nA, $\tau_{syn} = 2$ ms. In addition, two further variations of these models were used. In one case there was a static (DC) conductance with current fluctuations (typical parameter values: $E_{e} = 0$ mV, $E_{i} = -75$ mV, $g_{0e} = 0.0121$ $\mu$S, $g_{0i} = 0.0573$ $\mu$S, $\sigma_{syn} = 1.0$ nA, $\tau_{syn} = 2$ ms), and in another case there was a static (DC) current added to fluctuating conductances (typical parameter values: $I_{syn0} = 0$ nA, $E_{e} = 0$ mV, $E_{i} = -75$ mV, $\sigma_{v} = 0.012$ $\mu$S, $\sigma_{0} = 0.0284$ $\mu$S, $\tau_{v} = 10.49$ ms, $\tau_{i} = 2.728$ ms).

All simulations were performed using the NEURON simulation environment (Hines and Carnevale, 1997), running on DELL computers (DELL Computer Corporation, Round Rock, TX, USA) under the LINUX operating system.

**Data acquisition and analysis**

Records of spike-release events at the soma with a temporal precision of 0.1 ms from simulation lasting 60 s–100 s for each parameter set constitute the data basis for the present investigation. In most cases, the coefficient of variation $C_{v}$, defined by the ratio between the S.D. and mean of the interspike intervals (ISIs)

$$C_{v} = \frac{\text{S.D. of ISIs}}{\text{mean ISI}}$$ (4)

could be deduced with reliable statistics. The statistical validity of the results was addressed by running simulations of up to 1000 s duration for representative parameter sets in both simplified single-compartment models and detailed multi-compartment models. Analyzing the $C_{v}$ as a function of the simulation length showed that in all cases even at low firing rates (approximately 2 Hz), the $C_{v}$ values deduced from the spontaneous-discharge activity converged already after about 50 s to values obtained from marked extended simulation periods (see Fig. 1B).

For a Poisson process with refractory period $T_{R}$, the $C_{v}$ only depends on the mean ISI and $T_{R}$ (see e.g. Bugmann, 1995; Christoudoulou and Bugmann, 2001):

$$C_{v} = \frac{\text{mean ISI} - T_{R}}{\text{mean ISI}}$$ (5)

To investigate the impact of a specific model parameter, simulations were performed by changing the parameter in question while keeping all other fixed. In general, these changes were accompanied by changes in the firing rate, and allowed to depict the $C_{v}$ as a function of the mean ISI. For a given data set, the $C_{v}$ as a function of the mean ISI was fitted using Eq. 5 with $T_{R}$ as a free parameter. The $T_{R}$ which yielded the best fit can be viewed as characterizing the “effective” refractoriness of the Poisson trains produced by the model. This effective refractory period is expected to depend on both cellular properties (channel kinetics, membrane excitability, passive membrane properties) as well as the characteristics of the driving synaptic background activity and, thus, will be different for different data sets.

To investigate if the recorded spike trains resembled that of a Poisson-type firing process, the ISIs were analyzed for exponential distribution. To that end, ISI histograms (ISIH) were constructed and fitted by gamma distributions

$$\rho_{\text{ISI}}(T) = \frac{ar(T)\exp(-rT)}{qT!}$$ (6)

where $\rho_{\text{ISI}}(T)$ denotes the probability for occurrence of ISIs of length $T$ and $r$, $a$ and $q$ are parameters. In addition, autocorrelograms (displaying the probability $p$) of the output spike trains were constructed to show independence of the spike events.

In order to link the subthreshold dynamics to the firing activity, we investigated the $C_{v}$ and the mean ISI as a function of the following measure of proximity to firing threshold

$$\Delta = \frac{\sigma_{v}}{V_{i} - \bar{V}}$$ (7)

where $\sigma_{v}$ denotes the S.D. of membrane potential fluctuations, $\bar{V}$ the mean membrane potential and $V_{i}$ the mean firing threshold. In all cases, $\sigma_{v}$ and $\bar{V}$ were estimated from voltage traces from which spikes were removed. Unless otherwise stated, a mean firing threshold $V_{i}$ of $-50$ mV was used. The latter was estimated from the models and showed little ($\leq 1$ mV) variations for changes in the parameters considered here. $\Delta$ can be viewed as an “effective” inverse distance to firing threshold and will be referred to as “threshold accessibility”. Higher values of $\Delta$ correspond to an activity closer to threshold, either caused by a larger fluctuation amplitude $\sigma_{v}$ or a reduced distance between the average membrane potential $\bar{V}$ and the firing threshold $V_{i}$.

**RESULTS**

**High discharge variability in the detailed biophysical model**

We first used a compartmental model of a neocortical pyramidal neuron with active dendrites (see Experimental Procedures). This model was constrained from in vivo intracellular measurements of the impact of background activity (Paré et al., 1998; Destexhe and Paré, 1999). In agreement with in vivo recordings, the model displayed a depolarized $V_{m}$, high-frequency $V_{m}$ fluctuations and a marked decrease of input resistance (high-conductance
This result can be explained by pointing to the tight relation between an increase in the cell’s firing rate from 0.2 to 8 Hz; further conductances from 50% to 200% of the standard values (see Experimental Procedures; Fig. 1C, black dots). In both cases, the output spike train showed a high variability, quantified by a C_v between 0.8 and 1.0 (Fig. 1C). Increasing the quantal conductances from 50% to 200% of the standard values leads to increased firing rate fluctuations, quantified by the S.D. v_m, and the spiking response. The increase in the synaptic release frequency of both excitation and inhibition only minimally changed C_v between 3.5 and 4 mV. The membrane’s threshold accessibility Δ changed only minimally, with higher frequencies leading to a shift towards the firing threshold (increase in Δ) and shorter mean ISIs (see Fig. 2A, white dots). This indicates a stable balance between inhibition and excitation. In contrast, for changes in the quantal conductances, the C_v increased from 1.9 mV to about 4 mV, reflecting in a larger effective shift towards firing threshold in the investigated parameter range (Fig. 2A, black dots). In both cases we observed a saturation in the fluctuation amplitude, meaning that no further increase in C_v and, thus, the firing rate could be obtained even for very high synaptic conductances or release rates. Therefore a C_v for lower mean ISIs could not be deduced. However, fits of the available data show that in all cases the output spike trains resemble Poisson trains. This conclusion is further supported by ISIHs and autocorrelograms (Fig. 1D), which show that these spike trains are both exponentially distributed according to a gamma distribution and independent.

In a second series of simulations we addressed the effect of the statistical properties of background activity on the variability of discharges. To this end, we altered the correlation among the spatially distributed synapses. An increase in the correlation (between 0 and 0.3 Pearson’s correlation coefficient) led to a marked increase in C_v (between 1.5 and 6.5 mV for the standard model) and, thus, an effective shift towards firing threshold (Fig. 2B). This shift was accompanied by an increase in the firing rate (between 0 and 50 Hz for the standard model) for all investigated membrane excitabilities. Surprisingly, despite the differences in the threshold accessibility Δ, the output spike trains showed a high irregularity with C_v values between 0.8 and 1.0 (Figs. 3A and 2B) with exponentially distributed and independent ISIs (for a representative example see Fig. 1D), resembling the behavior expected from a Poisson process with refractory period. This suggests that, under physiological conditions, the appearance of highly irregular discharge patterns only weakly depends on the effective distance to firing threshold and, thus, the operating point in the subthreshold regime.

Interestingly, for comparable mean ISIs, the effective distance to firing threshold was smaller (larger Δ) for the standard setup (see Fig. 2B, right, black triangles). For higher excitability membranes, the higher membrane conductance shifts the average membrane potential towards the leak reversal potential. In addition, in more excitable membranes, smaller membrane potential fluctuations are needed to elicit spikes. Both effects together yield, compared to models with lower excitability, smaller Δ values for evoking spontaneous discharges at a given frequency (Fig. 2B, compare white dots with black triangles). On the other hand, if the membrane excitability is too low (Fig. 2B, black dots), the overall smaller membrane conductance brings the average membrane potential closer to threshold, whereas higher fluctuations are needed in order to elicit spikes at a given frequency. Both opposing effects yield an effective distance to threshold which was, for a given mean ISI, slightly larger (smaller Δ) than that for models with higher membrane excitability (Fig. 2B, compare black dots with black triangles).

We further investigated the impact of the membrane excitability on the effective distance to threshold and its role in modulating the variability of discharges by scaling the ion channel densities (see Experimental Procedures) for fixed background activity. An increase in the overall membrane excitability had a major impact on the average firing rate displayed by the cell, but not on the C_v (Fig. 3A). For a broad range of membrane excitabilities, between 80% and 200% around the standard values, the C_v was about 0.8 for firing rates up to 20 Hz (Fig. 3B). When the membrane excitability was less than 80% of the standard value, the firing rate was low (<0.2 Hz) for obtaining statistically reliable estimates of the C_v. As mentioned above, there was only a weak relation between the irregularity of the discharge and the threshold accessibility Δ and, thus, the membrane excitability (Fig. 2C, left). However, for increasing excitability, which was accompanied by a decrease in the mean ISI, Δ increased (Fig. 2C, right), indicating an effective shift towards firing threshold. For a further increase in the membrane excitability and mean firing rate, the opposite was seen, namely an increase in the distance (decrease in Δ, Fig. 2C, right). This suggests that for an intermediate excitability, the opposing effects of increased membrane potential fluctuations due to higher excitability and the shift towards the resting potential due to an increase in the overall membrane conductance at higher excitability counterbalance. The latter results in an effective shift towards the firing threshold but had no obvious impact on the discharge irregularity.

We also tested the dependence of the variability of the kinetics and the type of ion channels present in the mem-
Fig. 2. Characterization of discharge activity as a function of the threshold accessibility. The $C_v$ (left) and the mean ISI (right) are shown as functions of $\Delta$, defined as the ratio between the membrane potential fluctuation amplitude $\sigma_v$ and the distance between the firing threshold $V_T$ fixed to $-50$ mV in all cases, and the mean membrane potential $V$ (Eq. 7). (A) Results for different levels of background activity, obtained either by a change in the quantal synaptic conductances, or the release frequency at excitatory and inhibitory synaptic terminals (ratio was kept constant) around the standard parameter setup (see Experimental Procedures). The data correspond to the one shown in Fig. 1C. (B) Results for the high and low excitable as well as the standard model (see Experimental Procedures). The firing rate was changed by altering the correlation in the synaptic background activity. (C) Results for different levels of membrane excitability in the presence of correlated background activity (Pearson correlation coefficient of approximately 0.1).
Fig. 3. (Caption overleaf).
brane. As shown in Fig. 3C, the $C_V$ was always between 0.8 and 1.2 for spontaneous firing frequencies lower than 10 Hz. Fits to the $C_V$ as a function of the mean ISI (curves in Fig. 3C) as well as the ISI histograms (Fig. 4) along with

![Sep. 4. Typical ISIHs and autocorrelograms (insets) for various ion-channel settings (see Fig. 3C). (A) Voltage-dependent conductances including sodium $I_{Na}$ and delayed-rectifier potassium channels $I_{Kd}$ as well as A-type potassium channels $I_{KA}$ according to (Migliore et al., 1999). Fit: Eq. 6 with $q=2$, $r=0.073$ ms$^{-1}$, $a=2.012$. (B) Settings as in A with additional persistent sodium current $I_{NaP}$ (French et al., 1990; Huguenard and McCormick, 1992; McCormick and Huguenard, 1992). Fit: Eq. 6 with $q=2$, $r=0.069$ ms$^{-1}$, $a=1.744$. (C) $I_{Na}$, $I_{Kd}$ and $I_{M}$ (standard model, see Experimental Procedures) with additional $Ca^{2+}$-dependent potassium current (C-current) $I_{KCa}$ (Yamada et al., 1989) and a high-threshold $Ca^{2+}$-current (L-current) $I_{CaL}$ (McCormick and Huguenard, 1992). Fit: Eq. 6 with $q=3$, $r=0.079$ ms$^{-1}$, $a=1.869$. (D) $I_{Na}$, $I_{Kd}$ and $I_{M}$ (standard model, see Experimental Procedures) with additional persistent sodium current $I_{NaP}$. For more details see Experimental Procedures. All plots were obtained for a spontaneous activity at about 15 Hz firing rate (see star in Fig. 3C). Fit: Eq. 6 with $q=1$, $r=0.043$ ms$^{-1}$, $a=2.004$.

Fig. 3. High variability of discharge for various ion channel settings. (A) The $C_V$ as a function of the mean ISI for the high and low excitability as well as the standard model (see Experimental Procedures). The firing rate was changed by altering the correlation in the synaptic background activity. Fits: Eq. 5 with $T_R=11.1$ ms for high, $T_R=10.1$ ms for mid and $T_R=11.2$ ms for low excitability. Corresponding plots with the $C_V$ and mean ISI as functions of the threshold accessibility are shown in Fig. 2B. (B) $C_V$ versus mean ISI for different levels of membrane excitability in the presence of correlated background activity (Pearson correlation coefficient of approximately 0.1). Fit: Eq. 5 with $T_R=25.4$ ms. Corresponding plots showing the $C_V$ and mean ISI as functions of the threshold accessibility are presented in Fig. 2C. (C) The variability as a function of mean ISI for various ion channel settings and kinetics. White dots: voltage-dependent conductances including sodium $I_{Na}$ and delayed-rectifier potassium channels $I_{Kd}$ as well as A-type potassium channels $I_{KA}$, according to Migliore et al. (1999). Fit: Eq. 5 with $T_R=14.0$ ms. Black dots: $I_{NaP}$, $I_{Kd}$ and $I_{KA}$ conductances (Migliore et al., 1999) with additional persistent sodium current $I_{NaP}$ (French et al., 1990; Huguenard and McCormick, 1992; McCormick and Huguenard, 1992). Fit: Eq. 5 with $T_R=15.1$ ms. White triangles: $I_{Na}$, $I_{Kd}$ and $I_{M}$ (standard parameter setup, see Experimental Procedures) with additional $Ca^{2+}$-dependent potassium current (C-current) $I_{KCa}$ (Yamada et al., 1989) and a high-threshold $Ca^{2+}$-current (L-current) $I_{CaL}$ (McCormick and Huguenard, 1992). Fit: Eq. 5 with $T_R=26.4$ ms. Black triangles: $I_{Na}$, $I_{Kd}$ and $I_{M}$ (standard parameter setup, see Experimental Procedures) with additional persistent sodium current $I_{NaP}$. Fit: Eq. 5 with $T_R=10.2$ ms. The star refers to a spontaneous firing rate of about 15 Hz. For more details, see Experimental Procedures.
the flat autocorrelograms indicate that in all cases the spiking activity can be well described by a Poisson process with refractory period. Changes in the membrane time constant due to different ion channels directly translated into a change in the effective refractory period and, this way, a shift of the fitting curves. Representative ISIHs and autocorrelograms (spontaneous firing rate of 15 Hz, see star in Fig. 3C) for all cases are shown in Fig. 4. The histograms are well fit by gamma distributions, and the autocorrelograms are flat except for intervals shorter than the refractory period, supporting the conclusion that the output spike pattern resembles that of a nearly ideal Poisson process.

The threshold accessibility \( \Delta \) covered a broad range of values, with an increase in \( \Delta \) yielding a decrease in the mean ISIs (Fig. 5B, right) and a slight decrease in the CV (Fig. 5B, left) for all investigated active conductance settings. The latter decrease must be viewed as a direct result of the higher discharge rate, which is accompanied by the membrane operating closer to firing threshold (thus yielding larger \( \Delta \)), but also limits the irregularity of the discharge due to the impact of the refractory period. Taken together, the obtained results allow the conclusion that within the limits of the investigated parameter spectrum, the exact nature of the spike-generating mechanisms, the channel kinetics and channel distributions showed only minor effects on the Poisson structure of the spontaneous discharge and its firing variability.

Balanced inhibition and excitation was proposed as a possible mechanism for the high variability of cortical neuron discharges (Shadlen and Newsome, 1994; Bell et al., 1995; Feng and Brown, 1998, 1999, 2000; but see König et al., 1996). In order to test this hypothesis, we independently changed the excitatory and inhibitory release rates \( (v_{\text{exc}} \text{ and } v_{\text{inh}}, \text{ respectively}) \) within some range around their experimentally constrained mean values (see Experimental Procedures; Fig. 6A: 50%–300% and 70%–150% for excitation and inhibition, respectively). Such changes in the release rates shifted the average membrane potential relative to firing threshold due to changes in the mean inhibitory and excitatory synaptic conductances (see Fig. 5A). Here, diminishing inhibition or pronouncing excitation yielded a decrease in the effective distance to firing threshold (increase in \( \Delta \)), which was accompanied by an increase in the firing rate (increase in the mean ISI, see Fig. 5A, right) as well as a decrease in the discharge variability (Fig. 5A, left) due to the impact of the refractory period. However, for synaptic background activity determined to an increasing extent by excitation, \( \Delta \) decreased after passing a maximal value. Given the fact that changes in the firing rate at synaptic terminals directly impact the total membrane conductance, the appearance of such a maximum is comparable to the situation we encountered for variable membrane excitability (see Fig. 2C): although a higher activity at synaptic terminals increases the membrane-potential fluctuation amplitude (thus yielding an increase in \( \Delta \)), the increase in the overall membrane conductance accompanying a pronounced synaptic activity effectively diminishes \( V_m \) fluctuations and shifts the average membrane potential towards the leak-reversal potential (thus yielding an decrease in \( \Delta \)). Both effects counterbalance, resulting in a condition where the average membrane potential comes closest to the threshold for eliciting spikes.

Below 50 Hz, the variations of the CV as a function of the release rates (Fig. 6A) was entirely an effect of firing-rate modulation (Fig. 6B). There was a clear deviation from Poisson-type firing at higher firing frequencies (Fig. 6B, left star), which occurred when combining strong excitation with weak inhibition. In this case, the autocorrelograms deviated from a flat shape (Fig. 6C inset), which was due to a tendency to produce bursts of spikes (mostly doublets) with an ISI of about 8 ms, the typical membrane refractory period. This indicates that the cell, driven by the excitatory background, fires in a sustained way with high rates only constrained by the refractory period. On the other hand, for most of the cases, the firing rate was below 50 Hz and the corresponding CV was larger than 0.5 (Fig. 6B). In those cases, the ISIs were exponentially distributed (Fig. 6D). This behavior was present even for large disturbances in the release rates of excitation and inhibition, suggesting that the “balance” hypothesis does not apply to this model.

To further probe the impact of the balance between excitation and inhibition, we varied the excitatory release rate while keeping the inhibitory release rate constant. Fig. 7A shows that the CV was always high for all investigated unbalanced conditions. The Poisson structure of the output spike train depended on the statistics of the background activity. For low correlations, the deviation from the exponential distribution (Fig. 7B) was larger than for high correlations (Fig. 7C; both cases are for \( v_{\text{exc}} \) approximately 2 Hz, see star in Fig. 7A).

These results argue for an impact of the morphology on the output spike pattern: with increasing correlation, spatially distributed synaptic inputs become more synchronized, thus diminishing the role of the location of an individual synaptic stimulation. To investigate this further, we abstracted from the cell’s morphological structure by collapsing the dendritic tree into three compartments of same total membrane area, representing the perisomatic, proximal and distal regions (Fig. 8A, inset; see details in Experimental Procedures). Interestingly, the variability of the spiking output was slightly higher than in the detailed model (above 0.9 for firing rates below 50 Hz, see Fig. 8A), but showed only a minor dependence on the strength of the synaptic background activity. The latter was changed by scaling all quantal conductances with a common factor. A lower refractory period compared to the detailed model could be deduced from the fitting curves, indicating that membrane properties responsible for the refractory period do have a less effective impact in the case of a simplified morphology (compare with Fig. 3A). ISIH and autocorrelograms indicate only a slight deviation from the behavior expected from an ideal gamma distribution and independent Poisson process (Fig. 8B).
Fig. 5. Characterization of discharge activity as a function of the threshold accessibility. The $C_v$ (left) and the mean ISI (right) are shown as functions of $\Delta$, defined as the ratio between the membrane-potential fluctuation amplitude $\sigma_v$ and the distance between the firing threshold $V_T$ fixed to $-50$ mV in all cases (except for the setup with Ca$^{2+}$ currents in B, where $V_T = -44.5$ mV), and the mean membrane potential $\bar{V}$ (Eq. 7). (A) Results for different excitatory frequencies (scaled between 50% and 300% around the experimentally determined mean value of 1 Hz) and inhibitory frequencies $\nu_{inh}$ (scaled between 70% and 150% around the experimentally determined mean value of 5.5 Hz). (B) Results for various ion channel settings and kinetics. The data and setups correspond to Fig. 3C. (C) Results for various simplified models of synaptic background activity: a point-conductance model ($g_0$ + fluct $g$), a fluctuating current model ($i_0$ + fluct $i$), a model using a constant current with fluctuating conductance around zero mean ($i_0$ + fluct $g$) and a model using a constant conductance with fluctuating current around zero mean ($g_0$ + fluct $i$). Model parameters were for $g_0$ + fluct $g$: $g_0 = 0.0121 \mu S$, $g_{syn} = 0.0573 \mu S$, $0.003 \mu S \leq \sigma_{g_0} \leq 0.035 \mu S$, $0.006 \mu S \leq \sigma_{g_{syn}} \leq 0.07 \mu S$; for $i_0$ + fluct $i$: $-0.46$ nA $\leq i_{syn} \leq -0.34$ nA, $0.05$ nA $\leq \sigma_{i_{syn}} \leq 2.0$ nA; for $i_0$ + fluct $g$: $I_{syn} = -0.34$ nA, $0.002 \mu S \leq \sigma_{i_{syn}} \leq 0.03 \mu S$, $0 \mu S \leq \sigma_i \leq 0.08 \mu S$; for $g_0$ + fluct $i$: $g_{syn} = 0.0121 \mu S$, $g_0 = 0.0573 \mu S$, $0.5$ nA $\leq \sigma_{g_{syn}} \leq 2.5$ nA.
Simplified models of cortical neurons and discharge variability

The results obtained with the three-compartment model point towards a limit in reducing the morphology of the cell while keeping the same number of randomly releasing synapses. We therefore investigated another type of approach, in which background activity was represented by two fluctuating conductances, which parameters can be adjusted to reproduce the statistical properties of conductance and voltage fluctuations due to background activity (Destexhe et al., 2001). This point-conductance model consists of two conductances, excitatory and inhibitory, which vary according to a stochastic (Ornstein-Uhlenbeck) process (see Experimental Procedures). Including this point-conductance model in a single-compartment Hodgkin-Huxley neuron, or in real cortical neurons in vitro, can lead to in vivo-like characteristics, such as a high discharge variability (Destexhe et al., 2001).

Two of the parameters of this model, the mean excitatory and inhibitory conductances ($g_{e0}$ and $g_{i0}$), are best suited to probe the role of the excitatory–inhibitory balance on discharge variability. When varying these parameters, the spontaneous spike trains showed a high variability described by a $C_v$ above 0.6 (Fig. 9A) in the whole investigated range (between 20% and 250% around the standard values for both $g_{e0}$ and $g_{i0}$). This result compares to

Fig. 6. The impact of concurrent inhibition and excitation on the firing variability. (A) The $C_v$ as a function of excitatory frequency $v_{exc}$ (scaled between 50% and 300% around the experimentally constrained mean value of 1 Hz) and inhibitory frequency $v_{inh}$ (scaled between 70% and 150% around the experimentally determined mean value of 5.5 Hz). (B) $C_v$ values from A depicted as a function of the mean ISI. The data show a clear deviation from Poisson-like firing at very high firing rates (mean ISI=20 ms). The left star indicates deviation from Poisson type firing at firing frequencies above 100 Hz, the right star indicates a spontaneous firing rate of about 23 Hz. Fit: Eq. 5 with $T_r=17.8$ ms. Corresponding plots with the $C_v$ and mean ISI as functions of the threshold accessibility are shown in Fig. 5A. (C) ISIH and autocorrelograms (inset) for $v_{exc}=2.7$ Hz, $v_{inh}=4.4$ Hz and $v_{exc}=1.1$ Hz, $v_{inh}=4.4$ Hz (D). Whereas in the case of strong excitation the sustained high firing rate around 110 Hz (indicated by left star in B) causes a clear correlation between ISIs (peaks in the autocorrelogram in (C), inset) and, thus, a deviation from Poisson-type firing, at smaller firing rates around 23 Hz (indicated by right star in B) the cell returns to Poisson-like firing patterns (D). All results are from models with high membrane excitability (see Experimental Procedures) and correlated background activity (Pearson correlation coefficient approximately 0.1). Fits: Eq. 6 with $q=15$, $r=1.848$ ms$^{-1}$, $a=0.544$ (C) and $q=2$, $r=0.091$ ms$^{-1}$, $a=1.805$ (D).
Fig. 7. Irregular firing activity under unbalanced conditions. (A) The Cν as a function of the mean ISI for a fixed inhibitory synaptic release frequency of νinh=5.5 Hz but different excitatory frequencies for two background correlations (zero correlation and correlation corresponding to Pearson correlation coefficient approximately 0.1). Although the Cν under these strongly unbalanced conditions fits well with the behavior expected from Poisson-type activity, the ISIH and autocorrelograms for firing at around 50 Hz (star in A; B shows the results for zero correlation, C for high correlation) depart from the ideal Poisson-type behavior (stars). Both autocorrelograms indicate a preferred ISI of length equal to the effective membrane refractory periods (see fits in A). Fits: Eq. 5 with \( T_{ref} = 19.6 \) ms for \( c_{Pear} = 0.0 \) and \( T_{ref} = 15.6 \) ms for \( c_{Pear} = 0.1 \) (A) Eq. 6 with \( q=3, r=0.168 \) ms\(^{-1}\), \( a=1.046 \) (B) and \( q=4, r=0.277 \) ms\(^{-1}\), \( a=0.986 \) (C).
that shown in Fig. 6A. For a strong disturbance of the balance between excitation and inhibition, the cell showed a sustained firing at high rates with a preference for an ISI period of 8 ms, which can be viewed as an effective membrane refractory period (see Fig. 9D for an autocorrelogram and ISIH at 100 Hz; the deviation from the Poisson structure is indicated by the star). However, for a physiologically realistic spontaneous firing activity below 50 Hz, the ISIs are independent and follow an exponential distribution (see e.g. Fig. 10B).

A high spike variability with $C_v$ values above 0.85 was also obtained when the S.D. of the excitatory and inhibitory synaptic conductances, $\alpha_e$ and $\alpha_i$, or the corresponding time constants $\tau_e$ and $\tau_i$ were changed independently (Fig. 9B, C). However, in contrast to the aforementioned situation, the ISIHs and autocorrelograms had exponential distributions and independence of ISIs in all cases (data not shown). $\sigma_e$ and $\sigma_i$ are directly related to the correlation of excitatory and inhibitory synaptic inputs (Destexhe et al., 2001); therefore, these results are consistent with the detailed model. $\tau_e$ and $\tau_i$ are related to the decay time constants of excitatory and inhibitory synaptic conductances, suggesting a minor impact for these parameters as well.

To address the possible role of a high-conductance state in determining firing variability, we compared conductance-based and current-based models of background ac-

![Figure 8](image-url)

**Fig. 8.** Irregular firing activity in a three-compartment model (A) The $C_v$ versus mean ISI for two different strength of synaptic conductances ($a_{sc}$ denotes the scaling factor for synaptic conductances with respect to the standard parameter setup; see Experimental Procedures). The inset shows the morphology of the three-compartment model. The star indicates a spontaneous firing rate of about 15 Hz. Fits: Eq. 5 with $T_R=1.5$ ms for $a_{sc}=1$ and $T_R=1.7$ ms for $a_{sc}=2$. (B) ISIH and autocorrelogram (inset) at a firing rate of about 15 Hz (indicated by star in plot A) show that the spontaneous activity resembles nearly that of an Poisson process with small refractory period. Fit: Eq. 6 with $q=1$, $r=0.08$ ms$^{-1}$, $a=1.48$. 

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activity. In the fluctuating conductance model shown above, there was a static (DC) conductance ($g_{e0}$ and $g_{i0}$) with conductance fluctuations. We compared this model with a fluctuating current model, in which noise was injected in the membrane as a DC current with current fluctuations (see Experimental Procedures). We also used two further variations of this model. In one case there was a DC conductance with current fluctuations, and in another case there was a DC current added to fluctuating conductances (see details in Experimental Procedures). We found the highest $C_v$ for the two models having a DC conductance, despite the fact that these models allowed a more effective change in the mean firing rate (Fig. 10A). The $C_v$ was lower for the models with DC current, and lowest for the fluctuating-current model with values below 0.6 for firing frequencies above 5 Hz (Fig. 10A).

Marked differences between the discharge activity in models where synaptic background activity was described by conductance or current noise were also revealed when the $C_v$ was investigated as a function of the threshold accessibility (Fig. 5C, left). Here, for a given $\Delta$, the $C_v$ was lower for the fluctuating-current model and highest for models of high-conductance states (fluctuating-conductance model as well as model with DC conductance). Interestingly, in the latter case there was nearly no change in the $C_v$ for a broad range of $\Delta$. Contrary, in the fluctuating-current model, the $C_v$ dropped significantly for increasing $\Delta$, i.e. for a decreasing effective distance to firing threshold, as a direct result of the higher discharge rate closer to threshold and the large effective refractory period due to the large membrane time constant in the low-conductance state. In contrast, there was little difference between conductance-and current-based models in the dependence of the mean ISI on the threshold accessibility (Fig. 5C, right).

To test whether the low $C_v$ in the current-based models is due to the small noise time constant $\tau_{\text{syn}}$, we varied this parameter while adjusting the S.D. of the synaptic current noise ($\sigma_{\text{syn}}$) to yield a constant spontaneous firing frequency of about 15 Hz. $\tau_{\text{syn}}$ values between 2 and 5 ms, which cover the range of time constants of typical fast excitatory and inhibitory synapses, did not lead to an increase in the variability of the output spike train (e.g., $\tau_{\text{syn}}=4$ ms yielded $C_v=0.58$). Only for a large increase in the noise time constant we obtained high variability. At the upper limit of the simulated parameter range, $\tau_{\text{syn}}=20$ ms, the $C_v$ was 0.9.

Another difference between conductance-based and current-based models was also apparent in the ISIHs and autocorrelograms. In models with high conductance, the
Fig. 10. Evidence that high firing variability is linked to high-conductance states. (A) Comparison of various simplified models of synaptic background activity. In the point-conductance model \((g_0 + \text{fluct } g)\), excitatory and inhibitory conductances were changed around a mean according to an Ornstein-Uhlenbeck process, leading to a high \(C_v\) around unity. In the fluctuating current model \((i_0 + \text{fluct } i)\), random currents around a mean described by an Ornstein-Uhlenbeck process were injected into the cell, leading to a lower variability in the spontaneous discharge activity. In two other models, using a constant current with fluctuating conductance around zero mean \((i_0 + \text{fluct } g)\), and a constant conductance with fluctuating current around zero mean \((g_0 + \text{fluct } i)\), higher \(C_v\) values were obtained, showing that high conductance states account for the high discharge variability. The star indicates a spontaneous firing rate of about 17 Hz. For model parameters and corresponding plots showing the \(C_v\) and mean ISI as functions of the threshold accessibility, see Fig. 5C. (B and C) Typical ISIHs and autocorrelograms (insets) for spontaneous firing rates around 17 Hz (indicated by the star in A) for the fluctuating conductance (model parameters for B: \(g_{e0} = 0.0121 \mu S, \sigma_{e0} = 0.015 \mu S, \sigma_g = 0.015 \mu S\)) and fluctuating current model (model parameters for C: \(I_{syn0} = 0.44 \text{nA}, \sigma_{syn} = 0.030 \mu S, \tau_{syn} = 0.26 \text{ ms}\)). (D) ISIHs and autocorrelograms (insets) for spontaneous firing rates around 15 Hz for the fluctuating current model with long noise time constant \((I_{syn0} = 0.44 \text{nA}, \sigma_{syn} = 0.26 \text{nA}, \tau_{syn} = 20 \text{ ms})\). The \(C_v\) was 0.9, but there is a clear deviation from a Poisson process (indicated by the stars). Fits: Eq. 6 with \(q = 1, r = 0.033 \text{ ms}^{-1}, a = 3.332\) (B), \(q = 11, r = 0.108 \text{ ms}^{-1}, a = 3.515\) (C) and \(q = 3, r = 0.11 \text{ ms}^{-1}, a = 3.53\) (D).
output spike train clearly showed the behavior of a Poisson process, with gamma-distributed and independent ISIs, and a refractory period of about 8 ms (Fig. 10B). On the other hand, the fluctuating-current models deviate from this behavior (Fig. 10C, D). For small \( \tau_{\text{syn}} \) of 2 ms we observed a much larger depletion, pointing towards an effective refractory period of about 40 ms. The cell fired preferentially with an ISI around 90 ms (Fig. 10C). For large \( \tau_{\text{syn}} \) of 20 ms, both the ISIH and the autocorrelogram show a clear deviation from a gamma distribution (Fig. 10D, stars), indicating that the high \( C_V \) is not the result of an independent Poisson process, but rather that of an increased occurrence of spike bursts (doublets).

These findings may explain the difference in the \( C_V \) value in conductance-based and current-based models. The high-conductance state caused by the ongoing synaptic background activity leads to a marked decrease of the effective membrane time constant, which in turn translates the synaptic background activity into a cellular response with only minimal changes of its Poisson structure. On the other hand, current noise does not lead to a decrease in the membrane time constant, thus causing a severe filtering of short ISIs, which results in a decrease of the variability of the cellular response. In some cases, high \( C_V \) can be obtained in current-based models due to an increased occurrence of spike bursts.

**DISCUSSION**

In this paper we have investigated the spike statistics of spontaneously discharging non-bursting cortical neurons by using both detailed and simplified biophysical models of neocortical pyramidal neurons in the presence of excitatory and inhibitory synaptic background activity. The distributed background activity was adjusted to electrophysiological measurements in cat parietal cortex in vivo (Paré et al., 1998; Destexhe and Paré, 1999). We investigated how spike statistics and discharge variability are influenced by various parameters of the background activity. We found that neither its strength (quantal conductance, release rates) nor the balance between excitation and inhibition are stand-alone factors determining a high \( C_V \) for physiologically relevant firing rates. We rather found that highly variable \( C_V > 0.8 \) Poisson-distributed discharges can be evidenced as a stable qualitative property for a broad parameter range. Provided the neuron was within the limits of explored parameters (e.g. ion channel kinetics and distribution, membrane excitability, morphology) and parameter ranges, we failed to find significant effects of changes on the irregularity of spiking over that expected from a renewal process with a (effective) refractory period.

Intracellular measurements showed that the synaptic background activity present in vivo creates a high-conductance state with low input resistance (Matsumura et al., 1988; Baranyi et al., 1993; Contreras et al., 1996; Paré et al., 1998), and therefore of fast membrane time constant. Another prominent feature of in vivo recordings is the presence of high-amplitude (\( \alpha \) between 2 and 6 mV) voltage fluctuations (Azouz and Gray, 1999; Destexhe and Paré, 1999; Lampl et al., 1999). Including these characteristics in both detailed (morphologically reconstructed) and simplified (single-compartment) models, led to high \( C_V \) values for all models that had a high-conductance component. We therefore propose that the high-conductance state of cortical neurons is essential for maintaining an irregular firing activity in neurons receiving irregular synaptic inputs.

The main support for this proposition is provided by simplified models, in which it was possible to manipulate the excitatory and inhibitory conductances of background activity, and explore their parameter space in detail. We mapped the physiologically meaningful regions of parameters producing (a) highly irregular spontaneous discharges (\( C_V > 0.8 \)); (b) spontaneous firing rates between 5 and 20 Hz; (c) Poisson-distributed ISI intervals; (d) input resistance and voltage fluctuations consistent with in vivo estimates. We also compared fluctuating conductance models with fluctuating current models of equivalent voltage fluctuations, but with high input resistance. The latter type is often used to represent synaptic background activity in models (Bugmann et al., 1997; Sakai et al., 1999; Shinomoto et al., 1999; Svirskis and Rinzel, 2000) or in experiments (Holt et al., 1996; Hunter et al., 1998; Stevens and Zador, 1998). However we found here that current-based models lead to a more regular firing (\( C_V \) approximately 0.6 for 5–20 Hz firing frequencies) compared to conductance-based models. Intermediate models (high mean conductance with current fluctuations, or mean current with conductance fluctuations) displayed the highest \( C_V \) when the high-conductance component was present. This analysis indicates that the most robust way to obtain irregular firing consistent with in vivo estimates is to use neuron models in a high-conductance state.

These findings are not in conflict with recent results which outlined that the Ornstein-Uhlenbeck process does not reproduce cortical spiking statistics (Shinomoto et al., 1999). The latter model consisted in white noise injected directly as a current to the membrane, and is therefore equivalent to fluctuating current models. Injection of current as colored noise (Brunel et al., 2001), however, may lead to high \( C_V \), although not with Poisson statistics and only for particular values of the noise time constant (Fig. 10D).

Although the high discharge variability reported here is in accordance with previous conductance-based models, some of our conclusions depart from that of earlier contributions. In the single compartment Hodgkin-Huxley type model of cortical neurons investigated by Bell et al. (1995), a number of cellular and synaptic input parameters were identified which, if correctly combined, yield a balanced or “sensitive” neuronal state. Only in this state, which was characterized by a narrow parameter regime, the cell converted Poisson synaptic inputs into irregular output spike trains. This fine tuning causes the cell to operate close to the threshold for firing, which in turn leads to an input (noise)-driven cellular response. In addition, these authors showed an increase in the variability for stronger inputs. In this case, there was a net decrease of the membrane time
constant, which was identified as the cause of the irregularity (Bell et al., 1995). This is consistent with the necessity of a high-conductance state, as we report here, but the fine tuning required in the above study contrasts with the high robustness that we report here.

A single-compartment model with Hodgkin-Huxley type Na⁺ and K⁺ currents of hippocampal interneurons subject to Gaussian current noise and Poisson-distributed conductance noise was investigated in Tiesinga and José (1999). For Poisson-distributed inputs, these authors report a net increase in the $C_V$ for fixed mean ISIs for increasing noise variance, and a shift of the $C_V$ versus mean ISI to lower mean ISIs for increasing noise average. However, noise average and variance were quantified in terms of the net synaptic current, leaving a direct link between synaptic conductance and spiking statistics open.

The impact of correlation in the synaptic background on the neuronal response was also investigated by Salinas and Sejnowski (2000) using a conductance-based integrate-and-fire neuron. The authors conclude, in accordance with our findings, that the variability of the neuronal response in an intact microcircuit is mostly determined by the variability of its inputs. In addition, they showed that using a smaller time constant (similar to that found in a high-conductance state) leads to higher $C_V$ values, which is also in agreement with our findings. They also reported a decrease in the variability for an increase in the effective refractory period (relative to the time scale of changes in the postsynaptic conductances) caused by smaller synaptic time constants and larger maximal synaptic conductances. These results hold for a “balanced” model. However, a marked decrease of the $C_V$ was obtained when the “balanced” model was replaced by an “unbalanced” one. Although the sensitivity to the balance in the synaptic inputs was not investigated in detail, the model suggests a peak in the $C_V$ only for a narrow parameter range (see also Bell et al., 1995). The reported $C_V$ values of 1.5 for firing rates of 75 Hz were markedly higher than that found here, presumably because of the occurrence of bursts and significant deviations from the Poisson distribution, as we found here for high firing rates (Figs. 6C, 7, 8D).

Finally, to characterize how fluctuations impact on the discharge variability, we represented the $C_V$ as a function of a measure of “threshold accessibility” $\Delta$. This measure revealed differences between low-conductance and high-conductance states. In low-conductance states, the $C_V$ was dependent on $\Delta$ (as Fig. 5C, black dots), consistent with integrate-and-fire models (see e.g. Troyer and Miller, 1997). On the other hand, for high-conductance states, the $C_V$ was mostly independent of threshold accessibility, which is in overall agreement with our main finding, namely that the high discharge variability is highly robust to the details of the model, provided it is operating in a high-conductance state.

In conclusion, in all models examined here, including detailed compartmental models and simplified point neurons, the genesis of highly variable discharges ($C_V > 0.8$ and Poisson-distributed) by membrane-potential fluctuations was highly robust only in high-conductance states. This relation between high-conductance states and the variability of discharges should be further investigated by future models and dynamic-clamp experiments.

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